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THE EFFECTS OF VORTEX SHEDDING ON THE AERODYNAMIC PERFORMANCE OF AIRFOILS

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ABSTRACT

This is a comparative study involving various National Advisory Committee for Aeronautics (NACA) airfoils NACA 0012, NACA 0024, NACA 4412 and NACA 4424. The numerical simulation involves the air flow passing these models at high Reynolds number and various angles of attack. The study focuses mainly on the aerodynamic performance as represented by lift and drag coefficients. The Computational Fluid Dynamics simulation has been carried out in order to understand the phenomena of vortex shedding. The aim is to investigate the effects of vortex shedding on the aerodynamic performance of airfoils which has still been widely discussed over the years. The methodology has been validated; by using Spalart-Allmaras, K-epsilon and inviscid models, the first with the turbulence intensity and backflow of 0.10% contributed to relatively accurate result. The outcome reveals that there are few aspects that need to be considered in order to control or delay vortex shedding phenomena such as thickness-to-chord ratio, airfoil types, Reynolds number and whether or not it involves transient flow.

Keywords: vortex shedding, lift coefficient, drag coefficient, CFD, airfoil, spalart-allmaras, K-epsilon, inviscid flow, stall.

INTRODUCTION

Aerodynamic force fluctuations usually occur in fluid dynamic system due to vortex shedding. Vortex shedding is defined as fluctuating flow that takes place when fluid (air or water) passing over a bluff body at some definite velocities relying on the size and shape of the body. The process of vortex shedding phenomenon has been the subject of many researches. Even after years of efforts, it has still been a challenging issue in fluid dynamics.

Basically, the starting of vortex-sheet formation is an interaction between two separating shear layers. Once develops, vortex will continue to grow and move circularly from its connected shear layer. In some cases, it is shed downstream. The stability of vortex sheet was theorized by von Karman who argued that it is possible to stabilize vortex shedding if the vortices are shed alternately [1].

Due to the significance of vortex shedding, few related design modifications have been made. One of them customizing wing sections to improvement in aerodynamic performance. The presence of sharp edges on a bluff body is also important in improving the accuracy and consistency of the vortex shedding phenomenon. [2]. A body with sharp corner helps in clearing shedding of vortices, and results in high turndown ratio.

The vortex shedding process has been deeply studied by many researches [3], [4] and [5]. Eventually, it is an unsteady flow that takes place at certain range of velocities based on the geometry of the aerodynamic body.

There are two important properties involved in vortex shedding which is Reynolds number of the fluid flow and also the viscosity of the fluid passing over a bluff body. Although many studies on vortex shedding have been involving cylinders, the unsteady wake of fixed airfoils have also been experimentally studied in the

presence of shear layer separation and vortex shedding at both high and low Re [6], [7] and [8]. Note that the Re relates inertial forces to the viscous forces of fluid. For higher Re, viscous force less dominates the flow and therefore turbulence easily occurs, whereas for low Re, the viscous force dominates and laminar flow is thus developed. In the case of relatively high Re, the strong oscillations lead to the formation of a Karman Vortex Street.

There have been numerous intensive researches done in order to control vortex shedding. Such control is necessary due to, for instance, the shedding of vortices from alternate sides of bluff bodies is correlated with strong periodic transverse forces that can damage the body. Thus, suppression of vortex-induced vibrations is required. In some occasions, however, the shedding leads to reduction of drag, which is the desired feature in many applications [9].

In the case of circular cylinder, the vortex formation continues growing as long as it is fed by circulation from the separated shear layer [10]. Cases related to other shapes as aircraft wings contribute limited knowledge on vortex shedding.

This study investigates the effects of vortex shedding on the airfoil of aerodynamic performance by means of numerical simulation.

ANALYSIS

The effects of vortex shedding on the various types of airfoil were observed in ANSYS simulations. NACA 0012 was used to simulate different validation results with different model setups. Simulation results acquired were compared with the result from xfoil predictions. The purpose of running various turbulence models was to ensure the reliability of the method. Three different models used were viscous model of K-epsilon, Spalart-Allmaras and also inviscid model.

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Geometry and parameter setup

As shown in Figure-1, there are 4 different geometries of interest, particularly 2 symmetrical airfoils (i.e. NACA 0012 and NACA 0024) as well as 2 unsymmetrical airfoils (i.e. NACA 4412 and NACA 4424). The airfoils point data files that were utilized in this study were taken from [11]. The airfoils are different from each other in terms of their geometries and thickness. This can be observed from the NACA code itself. The last two digits indicate the thickness-to-chord ratio (in percent). Hence, NACA 0012 and NACA 4412 have 12% thickness-to-chord ratio. On the other hand, NACA 0024 and NACA 4424 both have the same 24% thickness-tochord ratio. Each airfoil was considered at different angles of attack in order to observe the variation in velocity contour and lift coefficient. Thus the stall angle of attack and the vortex shedding formation can be determined and detected, respectively. The problem set up and solution method with respect to turbulence models are shown in Table-1 and Table-2, respectively.

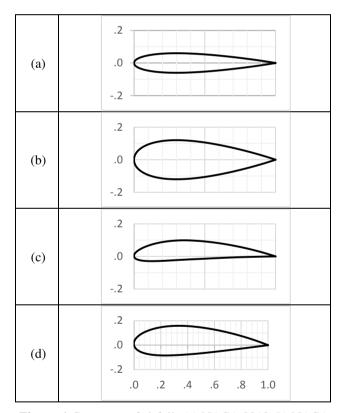


Figure-1. Geometry of airfoils (a) NACA 0012 (b) NACA 0024 (c) NACA 4412 (d) NACA 4424.

Meshing

The hexahedral mesh was used in the simulation. One of the mesh type advantages is that only less element needed for accurate solution. Thus, the solution time would be faster. Advanced size function was set as 'Proximity and Curvature' such that the value of maximum face size could be set manually. Such function

would also contribute to a much better mesh which adequately captures all the relevant features. The rest of the parameters were set as default. The mesh is shown in Figure-2.

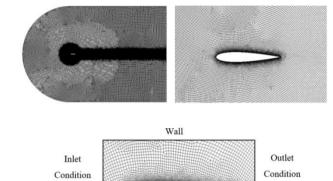


Figure-2. Grid applied for airfoil moving in unbounded flow.

RESULTS AND DISCUSSIONS

Method validation

V = 14.6073 m/s

The lift coefficient C_l profile changes over the time depending on the airfoil geometries. Various parameter combinations have been used in order to validate the data of NACA 0012. Both viscous and inviscid cases were considered. The turbulence models of K-epsilon and Spalart-Allmaras were used in the case of viscous flow. Figure-3 shows the curves of lift coefficient at various angles of attack for three models and against reference data. At low-to-moderate angles of attack, lift coefficient varies quite linearly with angle of attack, a. All the models have a good agreement with the reference data, except for inviscid model where C_l is almost linearly proportional to the angle of attack up to $\alpha = 20^{\circ}$. Thus, inviscid model is not suitable for reproducing the same data as the reference [12]. By way of contrast, Spalart-Allmaras turbulence model produces the same C_l behaviour as that of reference data such that the error is the lowest. This turbulence model is a simple one-equation model that models the kinematic eddy (turbulent) viscosity, is specifically designed for application, and it gives good results for boundary layers [13].

Once Spalart-Allmaras model was selected, it was tested using 3 different values of turbulence intensity and backflow which were 0.1%, 2.5%, and 5%.

Figure-4 shows the C_l profiles with respect to turbulence intensity and backflow. In general, the curves are similar to that of the reference data.

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Table-1. Parameters setup for three types of model.

	SE	ГИР		
Model	Spalart-Allmaras	K-epsilon	Inviscid	
Solver	Pressur	re Based	Density Based	
Velocity Formulation	Absolute			
Time	Steady			
2D Space		Planar		
	Mo	odel		
Energy	On	Off	On	
	Spalart-Allmaras (1 equation.)	K-epsilon	Inviscid	
Model	Spalart-Allmaras	K-epsilon model:		
	Production:	Standard Near wall treatment:		
	Strain/Vorticity based	Standard wall function		
	Material	Selection		
	Den	sity $(kg/m3) = Constant (1.$.225)	
Air	Cp (Specific Heat) (j/kg-k) = Constant (1006.43)		Cp (Specific Heat) (j/kg-k) = Constant (1006.43)	
	Thermal Conductivity			
	(w/m-k) = 0.0242 Viscosity (kg/m	l-s) = 1.7894e-05		
		Condition		
	1		Type : Wall	
	Type: Wall		Type. wan	
Free	Wall Motion = Stationary Wall			
Ticc	Shear Condition = No Slip			
	Roughness Height (m) = 0 (Constant) Roughness Constant = 0.5 (Constant)			
	Rouginiess Collsta			
	Type: Velocity-inlet			
	Velocity Specification Method = Magnitude and Direction Reference Frame = Absolute			
Inlet	Reference Frame = Absolute			
IIIIet	Velocity Magnitude (m/s) = 14.6073			
	Supersonic/Initial Gauge Pressure (pascal) = 0			
	X-Component of Flow Direction = 1 Y-Component of Flow Direction = 0			
		n = 0		
	Specification Method	elence: = Intensity and Length		
	Scale Turbulent Intensity (%) = 0.1%			
		$\sinh \text{ Scale (m)} = 1$		
	Turodicit Leng	Type: Pressure-outlet	<u> </u>	
	Gauge Pressure (Pascal) = 0			
	Backflow Direction Specification Method = Normal to Boundary			
Outlet	Turbulence:			
	Specification Method = Intensity and Length Scale			
	Backflow Turbulen	t Intensity (%) = 0.1		
	Backflow Turbulent	Length Scale $(m) = 1$		

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Table-2. Solution methods.

	SOLUTION	ON METHODS	SOLUTION METHODS				
Model	Spalart-Allmaras	K-epsilon	Inviscid				
	Pressure-Velocity Coup	ling	Formulation				
	Scheme : Coupled		Implicit				
			Flux type: Roe-FDS				
		Discretization					
Gradient		Least Squares Cell Based					
Flow			Second Order Upwind				
Pressure	Second (
Momentum	Second Orde	er Upwind					
Modified μ_t	Second Order Upwind						
Energy k	Second Order Upwind	Carand Onder Harrind					
1		Second Order Upwind Second Order Upwind					
3	Solut	ion Controls					
	Flow Courant Number :		Courant number : 5				
	Explicit Relaxation Factor		Courant number . 5				
	Momentum = 0.5	515.					
	Pressure = 0.5						
	Under-Relaxation Facto	rs ·					
	Density = 1	10.					
	Body Forces =1						
	Modified $\mu_t = 0.8$						
	$\mu_t = 1$						
	Energy = 1						
	N	Monitors					
		Create: Drag					
		Wall Zones = Part_1 & Part_2					
	Optio	Options : Print to Console and Plot					
_	Window = 2						
Residuals,		Average Over = 1					
Statistic and		Force Vector : X = 1, Y=0					
Force	***	Create: Lift					
Monitors		all Zones = Part_1 & Part_2					
	Options : Print to Console and Plot						
		Window = 3					
		Average Over = 1					
		Force Vector : X =0, Y=1					
Method	Solution Initialization						
Compute frm.	Standard Initialization Inlet						
Ref. Frame	Relative to Cell Zone						
Ref. I fame	G	Sauge Pressure (Pascal) = 0					
-	X Velocity (m/s) = 14.6073						
Initial Values							
	μ_t = 0.001610119	Y Velocity (m/s) = 0 $k = 3200598 \times 10^{-10}$	T(K) = 300				
	T(K) = 300	$\varepsilon_t = 5.725939e-6$	1 (11) 500				
	Run Calculation						
		Iterations = 3700					

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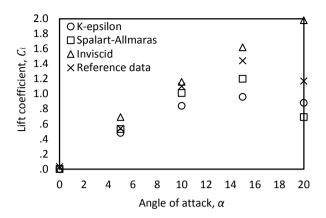


Figure-3. The lift coefficient, C_l profile.

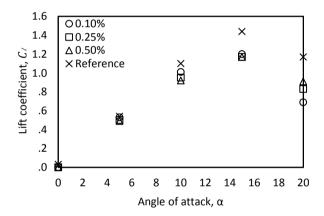


Figure-4. Comparison of C_l profiles with respect to turbulence intensity and backflow.

The lift coefficient profile in the case of 0.10% turbulence intensity is the closest to that of reference data; at $\alpha = 5^{\circ}$, $C_l = 0.53$ with only -1.85% error. The lift coefficient increases with angle of attack until it reaches the maximum.

The rest of the cases show greater errors compared to the previously mentioned. For instance, at $\alpha = 5^{\circ}$, C_l are 0.50 and 0.49 in the case of 2.5% and 5% turbulence intensity, respectively. This obviously shows greater errors.

Vortex shedding analysis

The velocity contour at stall angle of attack (i.e. on the left of the figure) and the vortex shedding (i.e. on the right of the figure) following the stall are shown in Figure-5. Since steady state solver was applied, the initially transient flow evolved into a steady flow. Theoretically, stall occurs when C_l is maximum, and vortex shedding begins right after the flow separation.

The velocity contours in the case of vortex shedding illustrate that the flow no longer follows the contour of the body (i.e. airfoil) at certain α . This shows that the flow has fully separated. Additionally, it can be

seen that the surrounding airstream velocity is very much higher than that of the airflow in the vicinity of each airfoil. Note that the bottom surface pressure distribution does not affect the separation. Two major consequences of the flow separations are a drastic loss of lift (i.e. stalling), and a major increase in pressure drag.

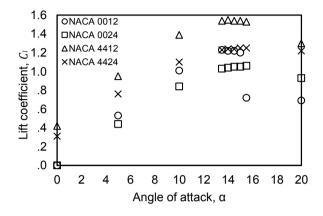


Figure-5. Diversity of lift coefficient for different airfoils at different angle of attack.

The lift coefficients for all four airfoils of interest are given in Figure-6. It is clear that at $\alpha = 0^{\circ}$, the lift coefficient is the same for all symmetrical airfoils where $C_l = 0$. Also, the C_l profile of NACA 0012 is greater than that of NACA 0024, in general.

The C_l profile of NACA 4412 is more 'curvy' in comparison to that of NACA 4424. Moreover, C₁ profile of NACA 4412 is higher.

In Table-3 the exact stall angle of attack and α when vortex shedding occurs are shown for each airfoil. The value of α is accurate to $\pm 0.5^{\circ}$. Since both symmetrical airfoils have less thickness, the stall occurs relatively earlier. Obviously, the location where vortex shedding first occurs is the trailing edge. Such occurrence vanishes when the steady state condition is achieved.

Vortex shedding was captured at $\alpha = 16^{\circ}$, 2.5° after stall for NACA 0012. Another symmetric airfoil case reveals that vortex shedding was captured $\alpha = 22^{\circ}$; the formation is delayed for 6.5° after stall. As far as unsymmetrical airfoils are concerned, the formation is captured at $\alpha = 18^{\circ}$ and 20° for NACA 4412 and NACA 4424, respectively.

Both the lift and drag coefficient values when the vortex shedding takes place are presented in Table-4. The lift-to-drag ratio l/d = 10.7 for NACA 0012 in comparison to l/d = 5.12 for NACA 0024 (i.e. differ by approximately 50%). Meanwhile, l/d = 11.83 and l/d = 11.09 for both unsymmetrical airfoils NACA 4412 and NACA 4424, respectively. This shows that an airfoil with less thickness gives greater *l/d* even in the presence of vortex shedding.



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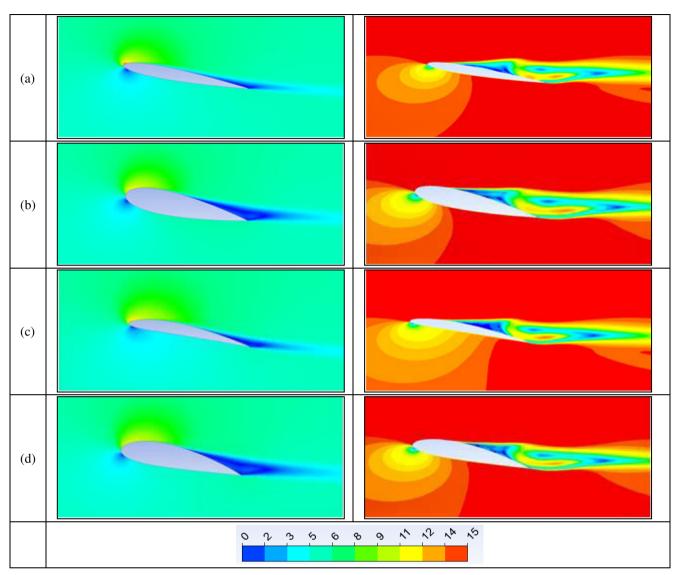


Figure-6. Velocity contour in m/s at α (a) NACA 0012 (b) NACA 0024 (c) NACA 4412 (d) NACA 4424.

Table-3. Occurrence of vortex shedding and stall.

Airfoils	Angle of attack, α		Location of first
	Stall	Vortex shedding	vortex shedding occurrence
NACA 0012	13.5°	16°	Trailing edge
NACA 0024	15.5°	22°	Trailing edge
NACA 4412	14°	18°	Trailing edge
NACA 4424	15°	20°	Trailing edge

CONCLUSIONS

The main purpose of this study is to investigate the effects of vortex shedding on the airfoil aerodynamic

performance as represented by lift and drag coefficients. NACA airfoils were selected based on their thickness and whether or not they are symmetrical. The approach of

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Table-4. The relationship between the angle of attack when vortex shedding occurs and aerodynamic performance.

Type of NACA	Lift coefficient, C_l	Drag coefficient, C_d
NACA 0012	1.07	0.10
NACA 0024	0.87	0.17
NACA 4412	1.42	0.12
NACA 4424	1.22	0.11

numerical simulation with steady state solver has been taken. The Reynolds number of 1,000,000 with uniform velocity of 14.6073 m/s were adopted.

Lift-to-drag ratio coefficient is higher for unsymmetrical airfoil than that for symmetrical airfoil even when vortex shedding presents. By increasing the angle of attack a, the energy of vortex shedding significantly appears. High \(\frac{1}{2} d \) is typically one of major goals in designing wings. Since an aircraft's required lift depends on its weight, delivering lift with lower drag directly leads to smaller fuel consumption, and better climb performance and glide ratio.

The way vortex shedding influences airfoils' performance can be seen in the context of airfoil geometries as their thickness, positions of both camber and chord lines, and whether or not they are symmetrical. Although vortex shedding cannot be totally prevented, it can be minimised or delayed.

The extension of this study would be, for example, the study of factors that influence the vortex shedding, as in the case of airfoils experiencing ground effects [14], [15], as well as the airfoils in compressible and flows [16], [17].

Important aspects to consider in vortex shedding control mechanism

The study provides information on how and when vortex shedding occurs. First, thickness to chord ratio of an airfoil has been examined; vortex shedding in the case of airfoil with lower thickness-to-chord ratio takes place earlier.

Another aspect is regarding the types of airfoil. It is understood that vortex shedding happens after stall. Since stall happens relatively early for unsymmetric type of airfoils, so does the vortex shedding.

Moreover, smaller Re tends to delay vortex shedding. Thus the parameter could be taken into account in the control mechanism of vortex shedding.

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